

Modeling the observed spectra and light curves of synchro-curvature emission of pulsars

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Motivation

Explain the high-energy spectra and light curves of pulsars, in gamma-rays and X-rays.

The model adopts an effective, but physically-based approach, versatile enough to:

- Fit the entire population of pulsars
- Test different emission regions
- Add/test new physics

FFE / PICS
/ FIDO

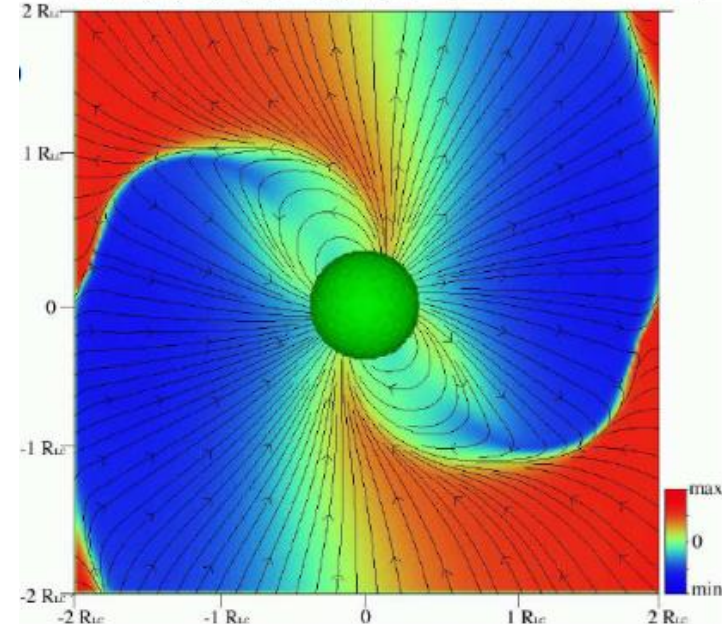
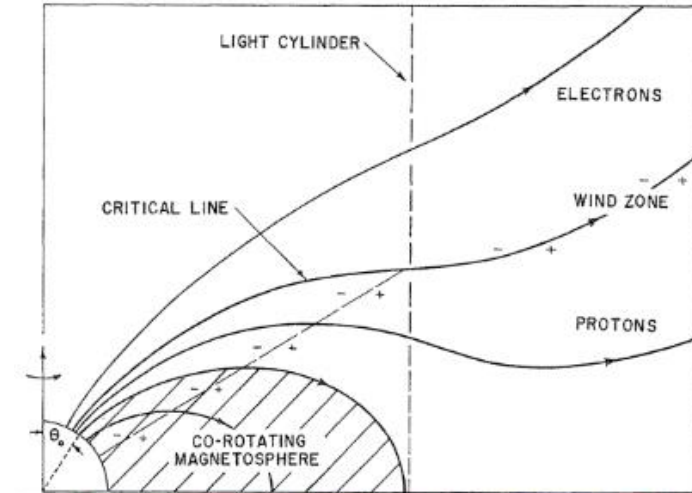
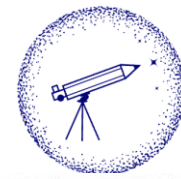
These
works

Phenom.
models

Versatility

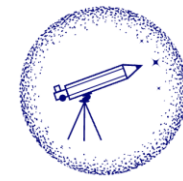
Computational cost

Physics

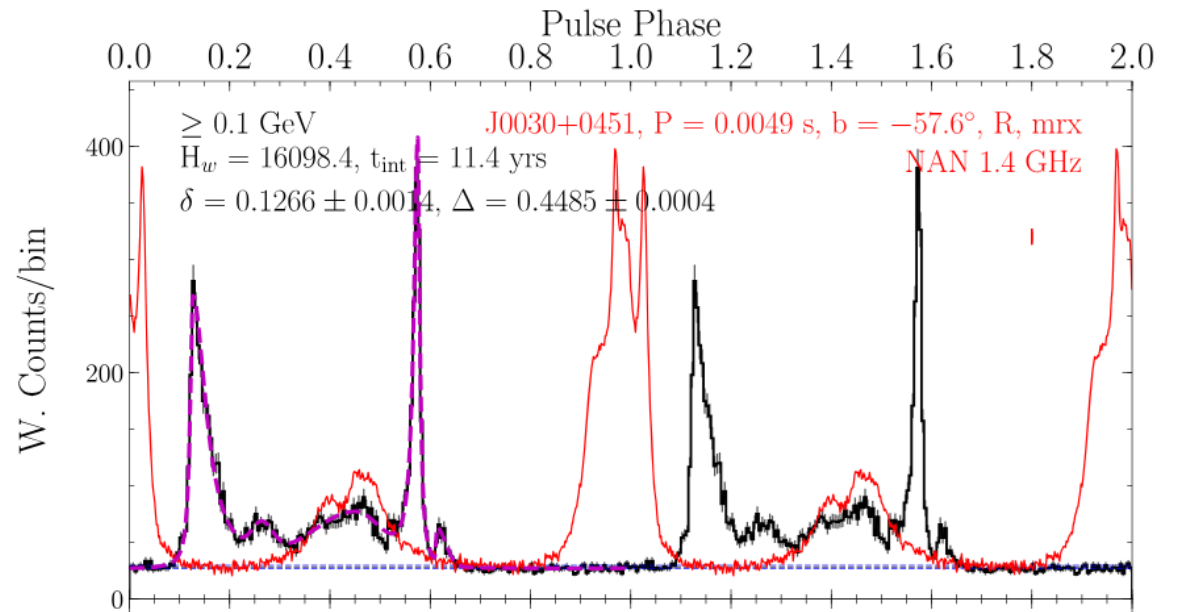
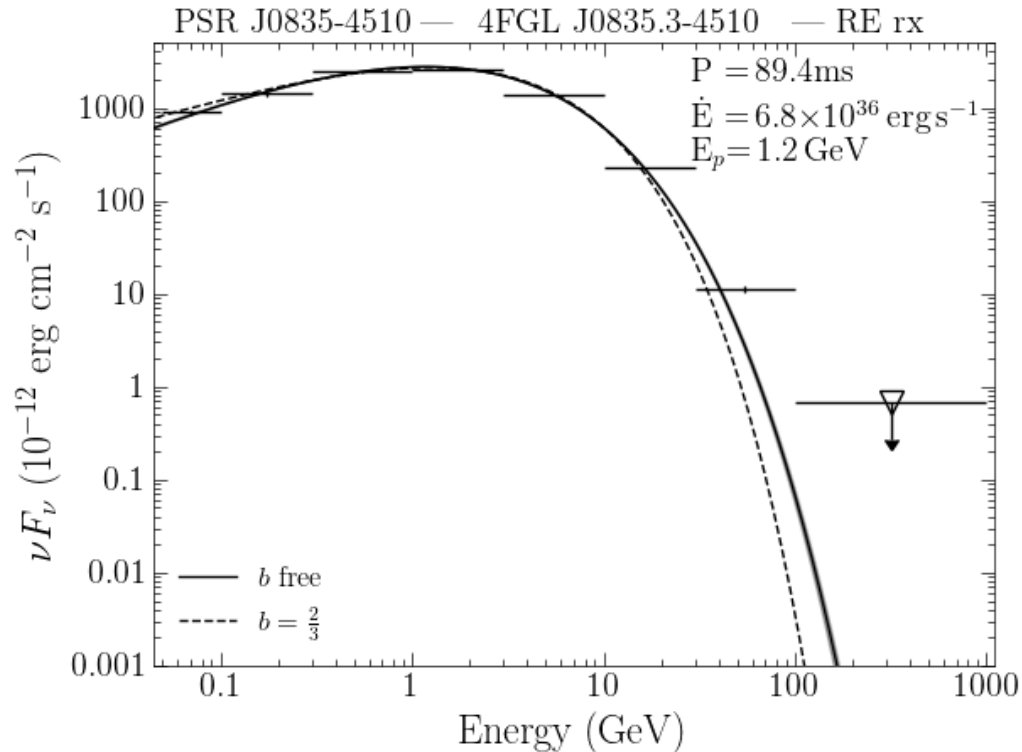


[Goldreich &
Julian (1969);
Spitkovsky
(2006)]

Motivation



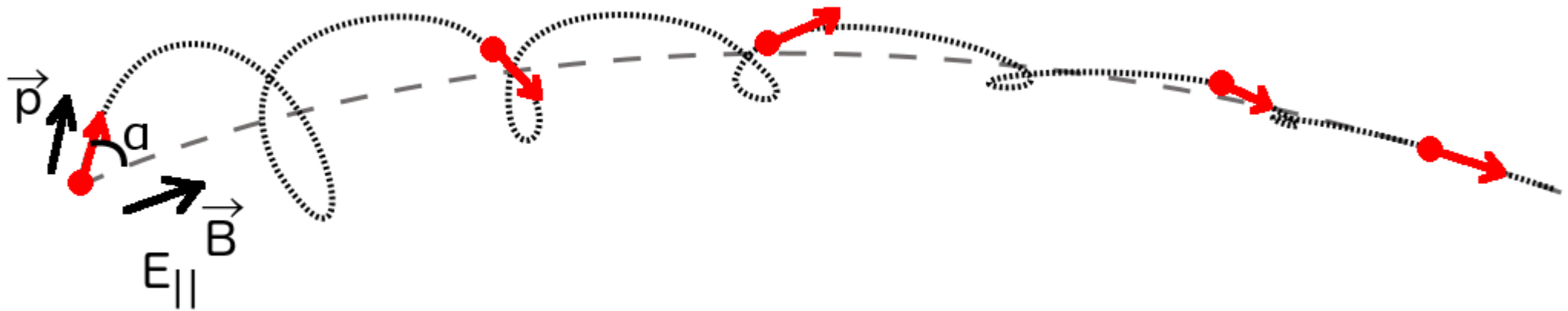
- Models aim to fit data of the gamma-ray emitting pulsars released in the 2PC and 3PC (in the latter, fewer spectral bins!)



- Our goal is to reproduce observational data with a model that contains simple but realistic physics and is computationally affordable

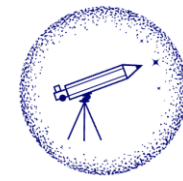
Spectral model: particle dynamics

- We follow the dynamics of the emitting particles, ruled by electric acceleration and synchro-curvature losses and with two free parameters involved: $E_{||}$, b
- Solving the equation of motion gives the evolution of the relativistic momentum and of the Lorentz factor Γ and pitch angle α



[Viganò et al. 2015, MNRAS, 447, 1164]

Spectral model: particle dynamics



- The equation of motion of charged particles balances electric acceleration and synchro-curvature losses

$$\frac{d\vec{p}}{dt} = ZeE_{\parallel}\hat{b} - (P_{sc}/v)\hat{p} \quad \vec{p} = \Gamma m\vec{v} \quad [\text{Viganò et al. (2014)}]$$

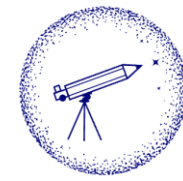
- We solve it numerically, considering separately the components parallel and perpendicular to the trajectory

$$\frac{d(p \cos \alpha)}{dt} = ZeE_{\parallel} - \frac{P_{sc}}{v} \cos \alpha \quad \frac{d(p \sin \alpha)}{dt} = - \frac{P_{sc}}{v} \sin \alpha$$

- Local magnetic field strength and curvature radius are parametrized in an effective way:

$$B = B_{\star} \left(\frac{R_{\star}}{x} \right)^b \quad r_c = R_{lc} \left(\frac{x}{R_{lc}} \right)^{\eta}$$

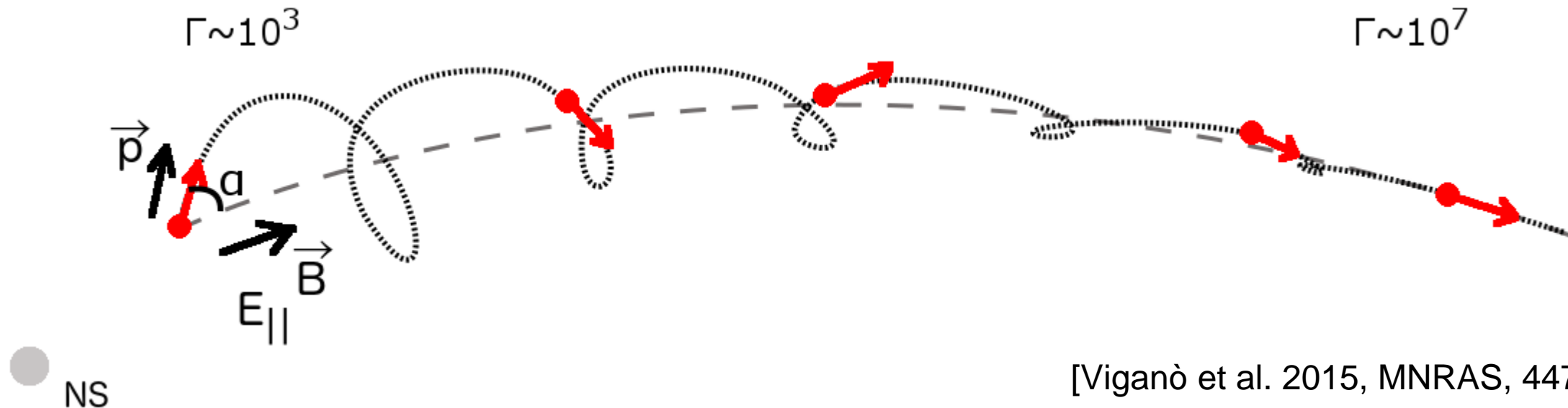
Spectral model: particle dynamics



- Power of synchrotron and curvature radiations depend on the Lorentz factor Γ of particles differently

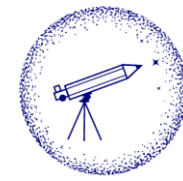
$$P_{syn} = \frac{2}{3} \frac{(Ze)^4 B^2 (\Gamma^2 - 1) \sin^2 \alpha}{m^2 c^3}$$

$$P_c = \frac{2}{3} \frac{(Ze)^2 c \Gamma^4}{r_c^2}$$



[Viganò et al. 2015, MNRAS, 447, 1164]

Spectral model: synchro-curvature radiation



- Single-particle synchro-curvature power spectra:

$$\frac{dP_{sc}}{dE} = \frac{\sqrt{3}(Ze)^2\Gamma y}{4\pi\hbar r_{eff}} [(1+z)F(y) - (1-z)K_{2/3}(y)]$$

- Synchro-curvature power:
$$P_{sc} = \frac{2(Ze)^2\Gamma^4 c}{3r_c^2} g_r$$

[Cheng & Zhang (1996),
Viganò et al. (2014)]

$$r_{eff} = \frac{r_c}{\cos^2 \alpha} \left(1 + \xi + \frac{r_{gyr}}{r_c} \right)^{-1}$$

$$z = (Q_2 r_{eff})^{-2}$$

$$y = \frac{E}{E_c}$$

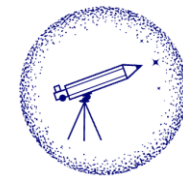
$$Q_2^2 = \frac{\cos^4 \alpha}{r_c^2} \left[1 + 3\xi + \xi^2 + \frac{r_{gyr}}{r_c} \right]$$

$$E_c = \frac{3}{2} \hbar c Q_2 \Gamma^3$$

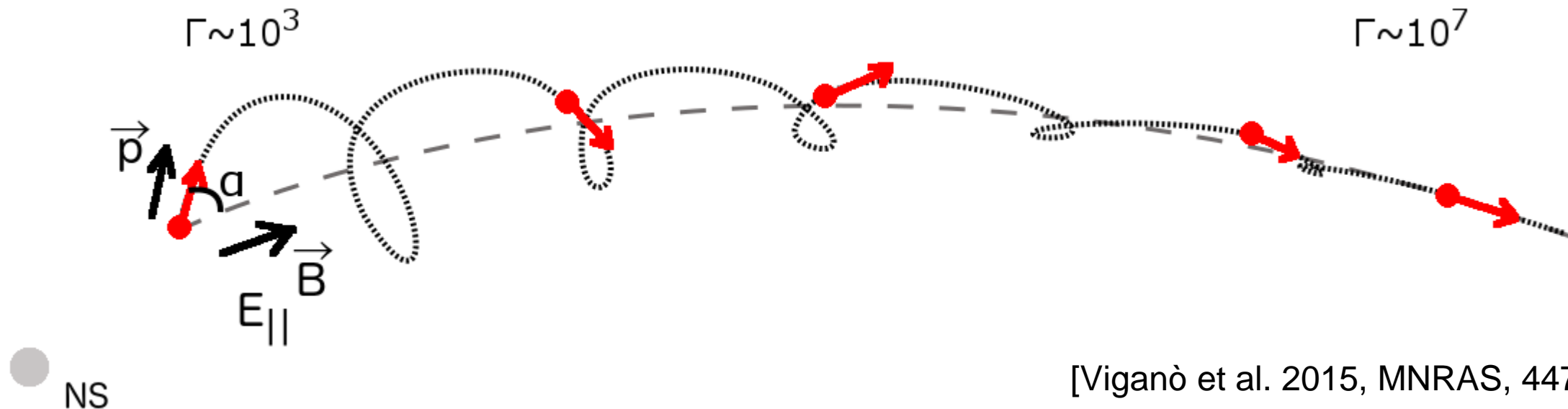
$$F(y) = \int_y^\infty K_{5/3}(y') dy'$$

$$\xi = \frac{r_c}{r_{gyr}} \frac{\sin^2 \alpha}{\cos^2 \alpha}$$

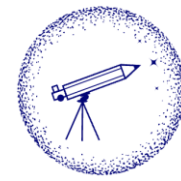
Spectral model: emission



- Synchro-curvature formulae gives the emission of the particles all along the trajectory, which convolved with an effective particle distribution gives the total radiation from the emission region
- We produce theoretical spectra with just three free parameter ($E_{||}$, b , x_0) and a normalization factor



[Viganò et al. 2015, MNRAS, 447, 1164]



- Convoluting the single-particle power spectra with an effective particle distribution,

$$\frac{dN}{dx} = N_0 \frac{e^{-(x-x_{min})/x_0}}{x_0(1 - e^{-(x_{max}-x_{min})/x_0})}$$

We obtain the total emission from the region:

$$\frac{dP_{tot}}{dE} = \int_{x_{min}}^{x_{max}} \frac{dP_{sc}}{dE} \frac{dN}{dx} dx$$

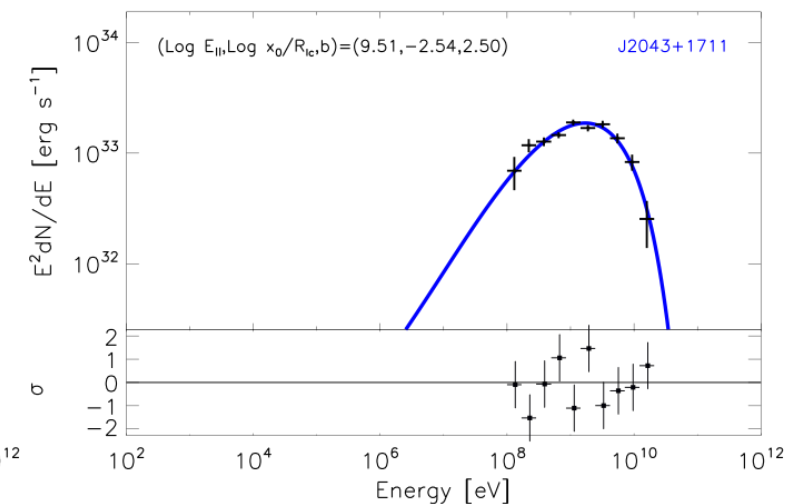
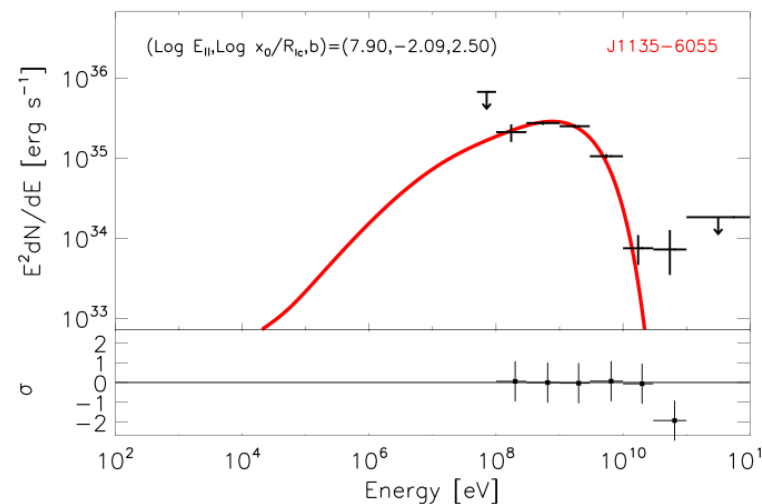
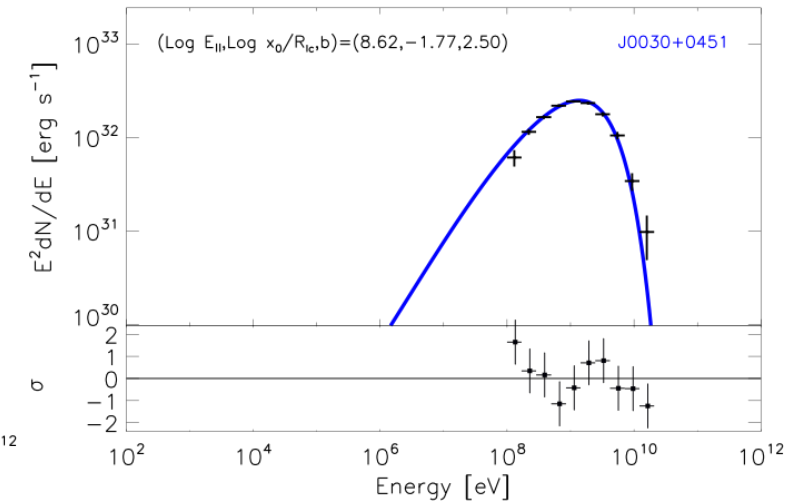
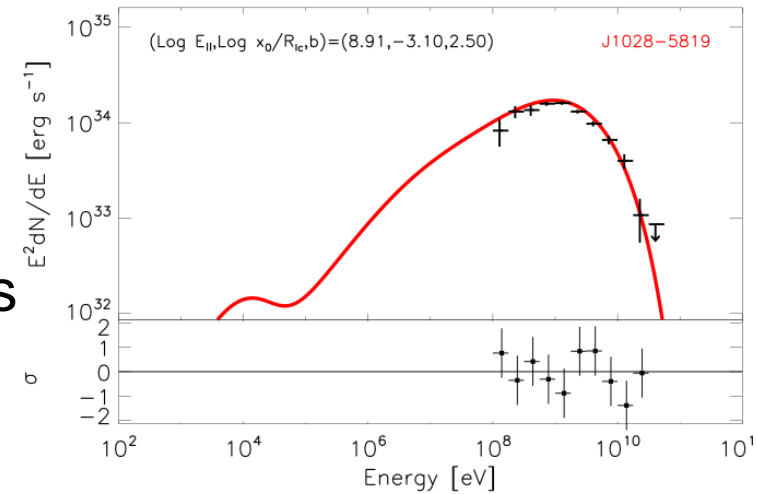
[Cheng & Zhang (1996),
Viganò et al. (2014)]

Spectral fitting to gamma rays

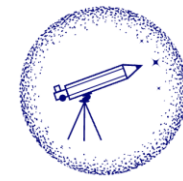
[Viganò et al. 2015, MNRAS, 453, 2599; Íñiguez-Pascual et al. (in prep.)]

- The model successfully fitted the 117 gamma ray-pulsars on the 2PC (Abdo et al. (2013)) and the ~300 gamma-ray pulsars on the 3PC (Smith et al. (2023))

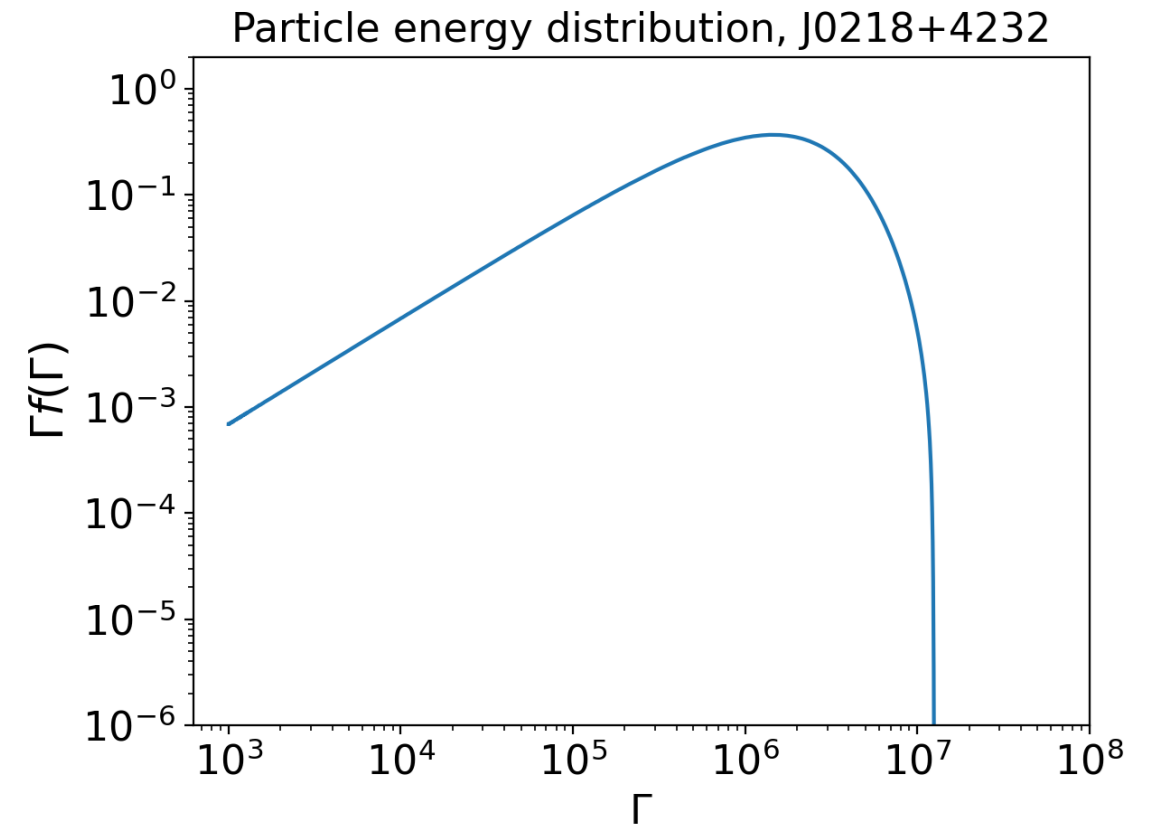
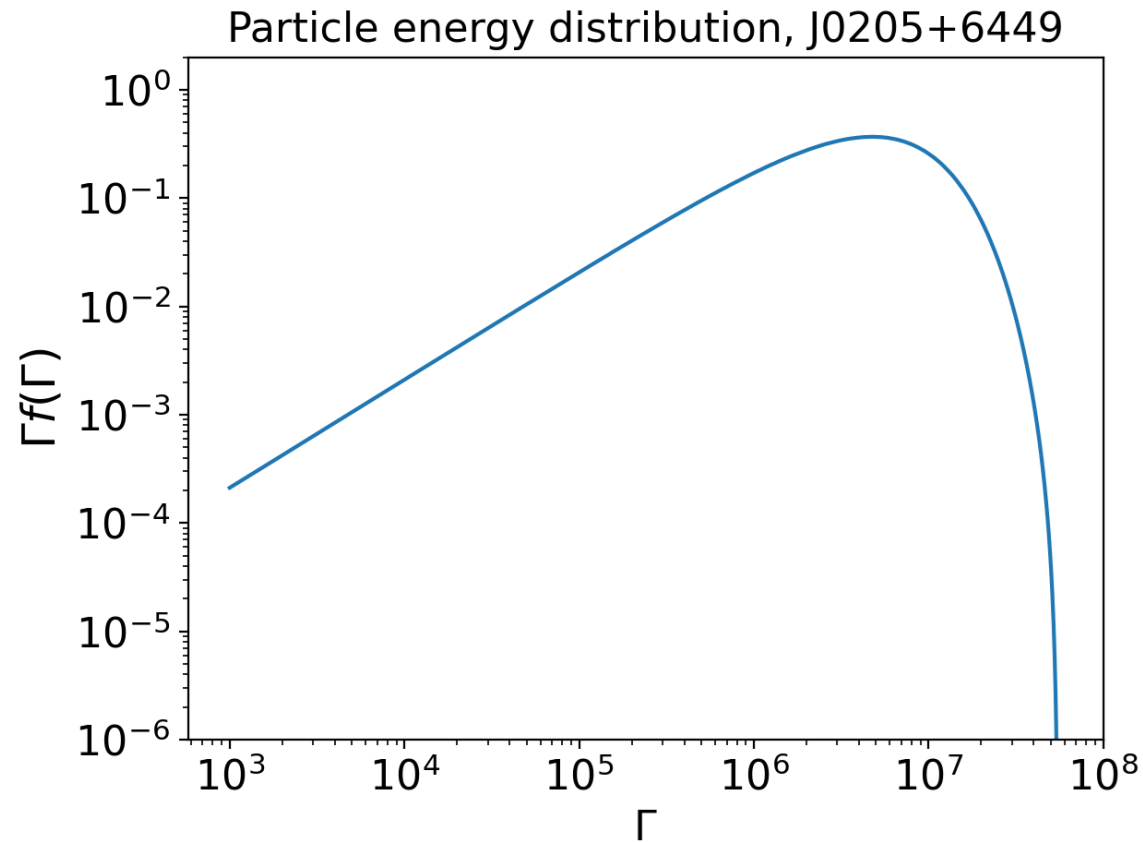
- A relevant synchrotron contribution is needed to match gamma-ray spectra, implying that synchro-curvature radiation is an appropriate mechanism to explain the emission from these objects



Spectral model: particle energy distribution

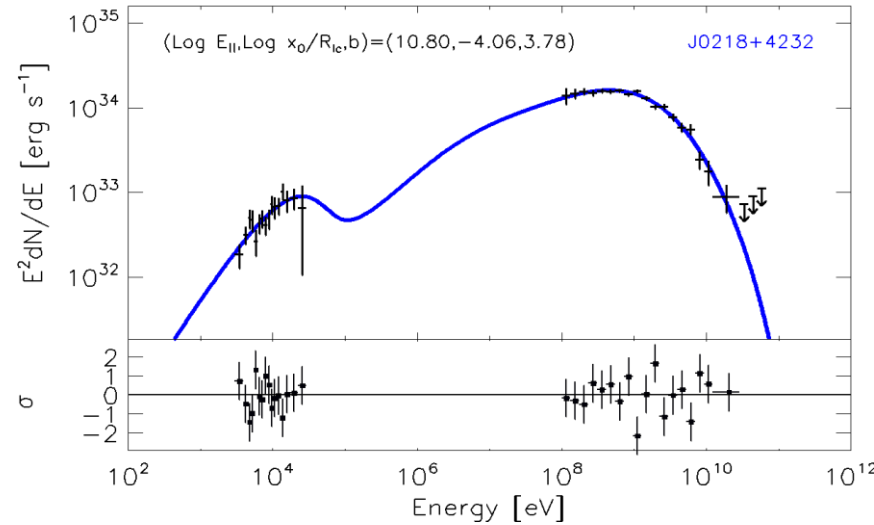
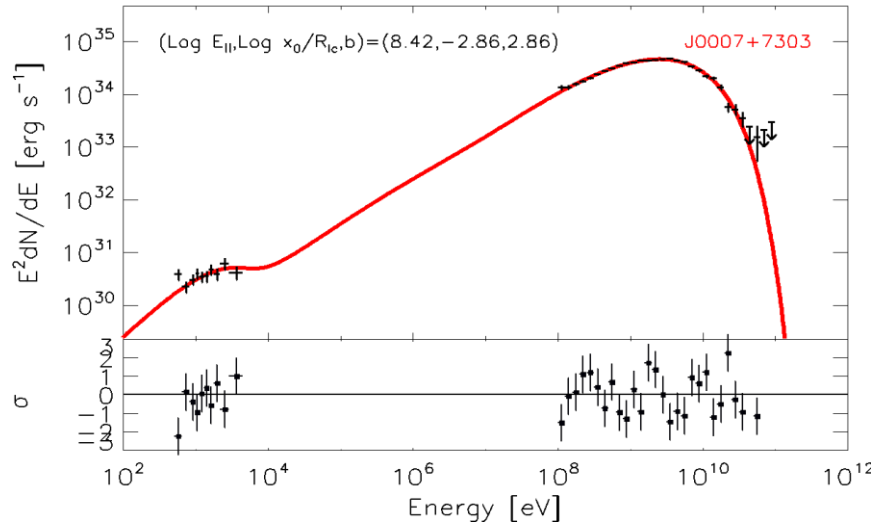
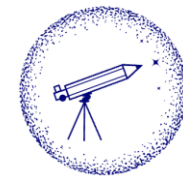


- Lorentz factors Γ typically range from 10^3 to 10^7 .

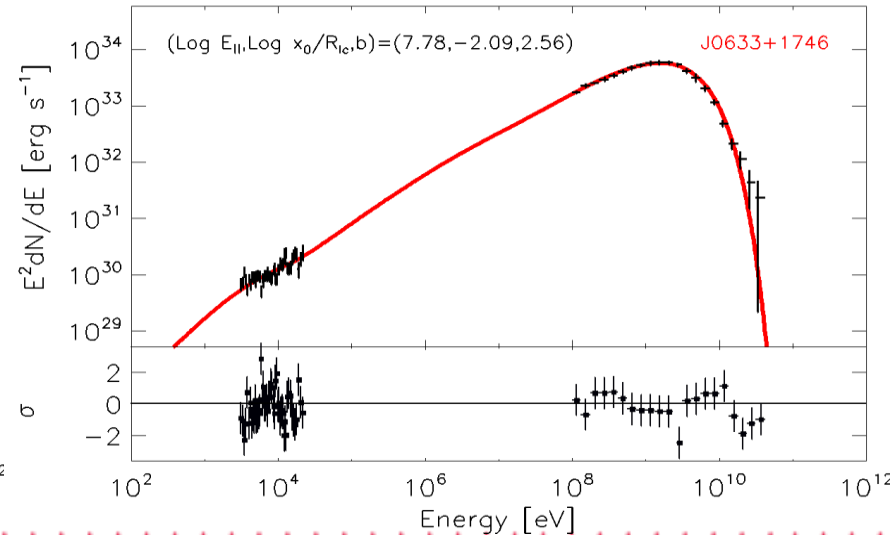
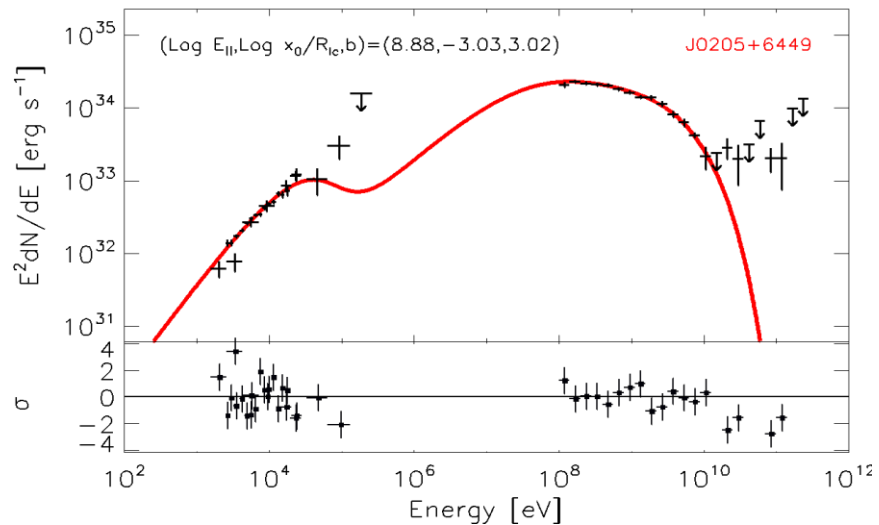


Spectral fitting

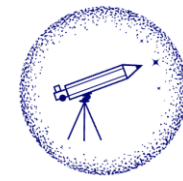
[Torres 2018, Nat. Astron., 2, 247;
Torres et al. 2019, MNRAS, 489, 5494;
Coti Zelati et al. 2020, MNRAS, 492, 1025;
Íñiguez-Pascual, Viganò & Torres 2022, MNRAS, 516, 2475]



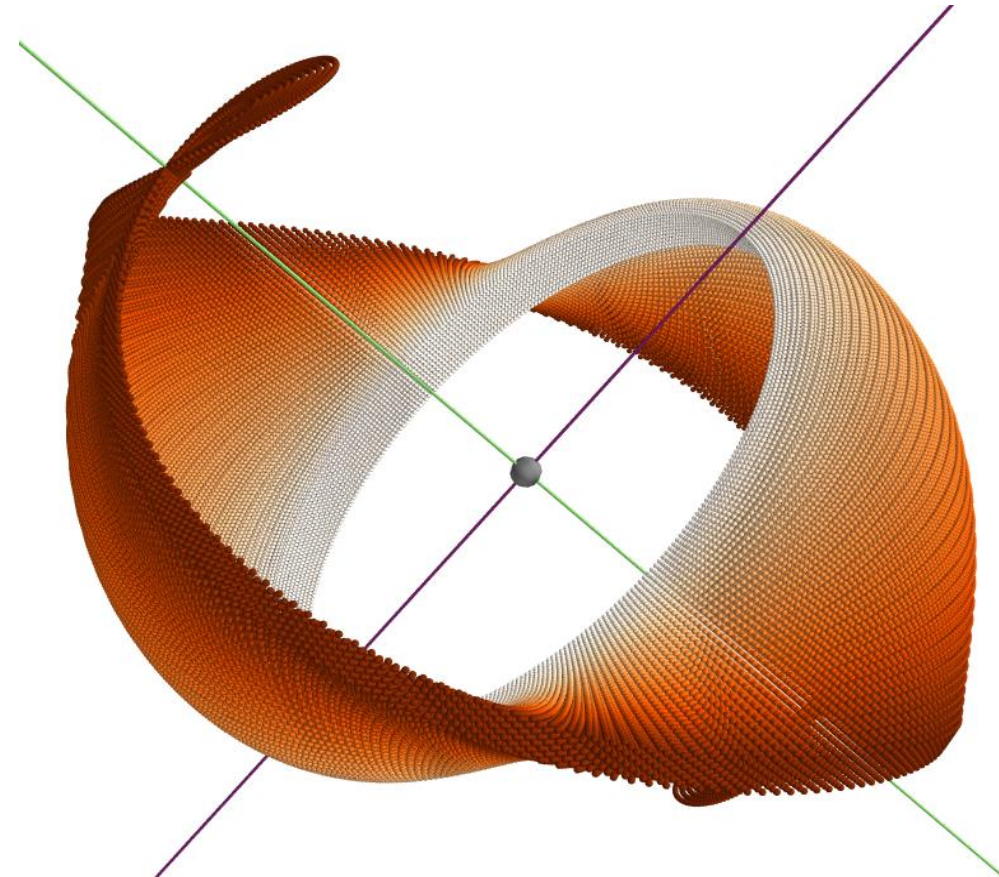
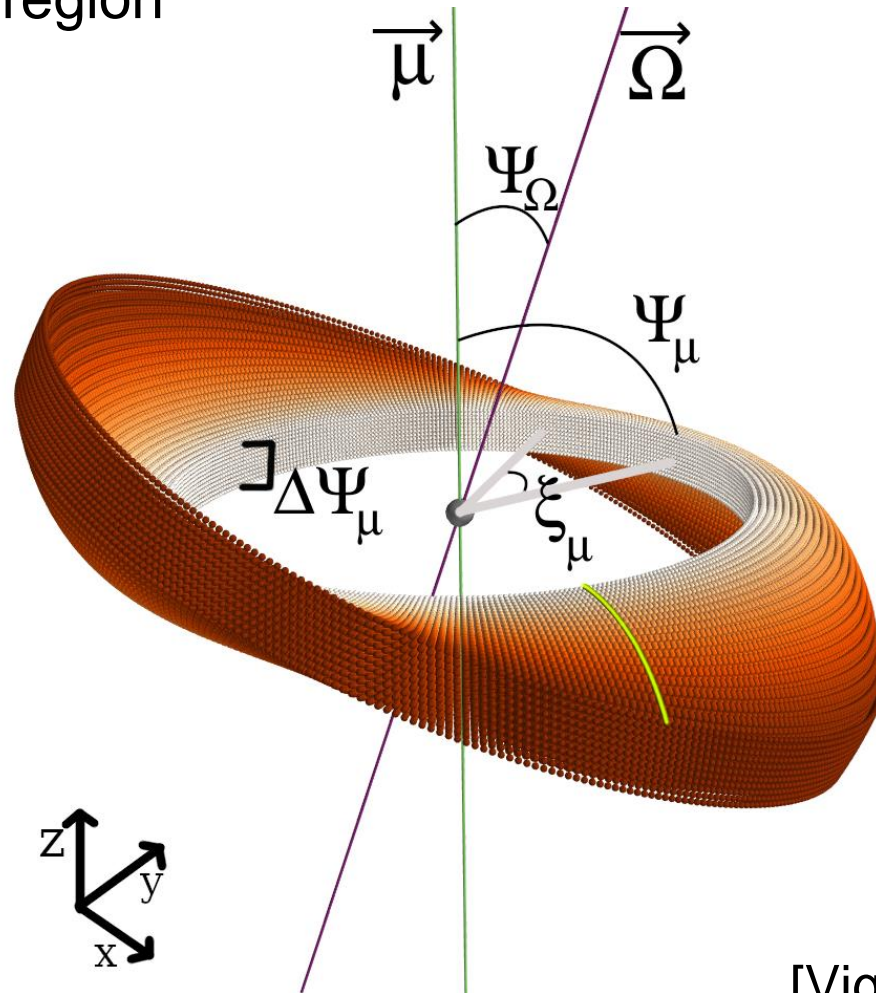
Extension to fit X-rays
for the ~40 X+ γ pulsars



Geometrical model

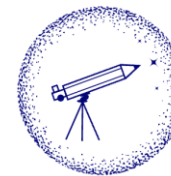


- The inclination angle and the meridional extent define the geometry of the emission region

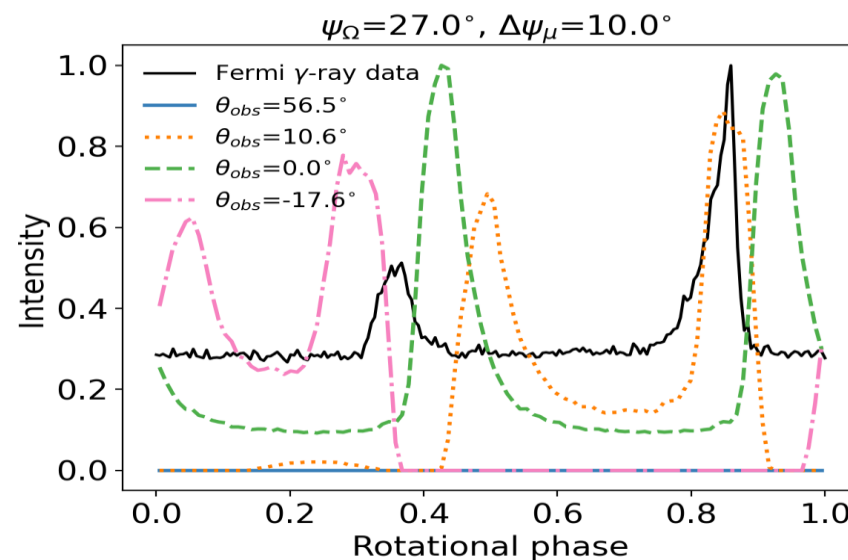
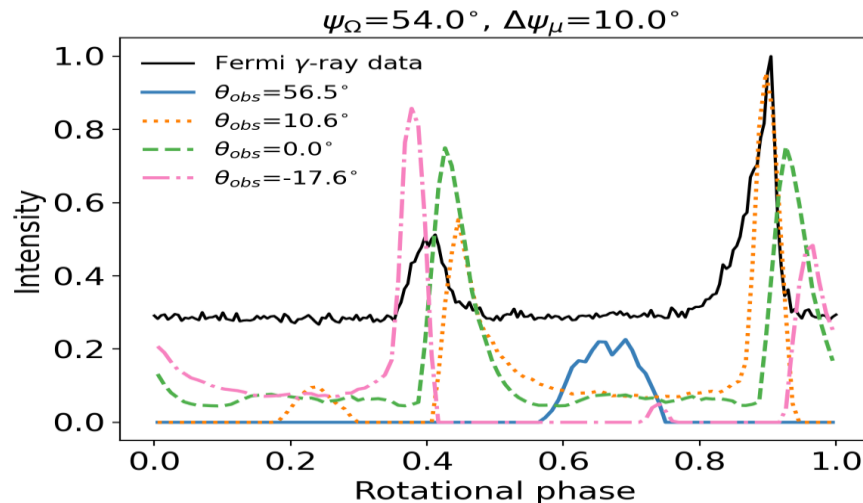
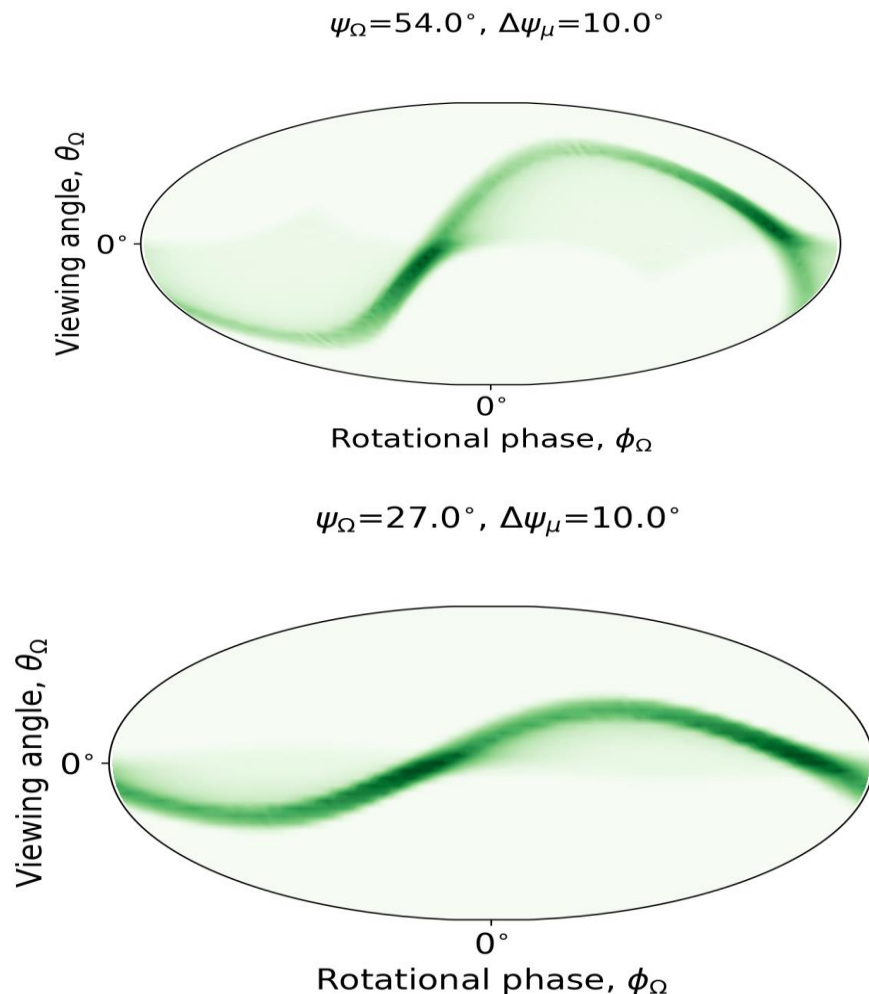


[Viganò & Torres (2019); Íñiguez-Pascual, Torres & Viganò (2024)]

Emission maps and light curves



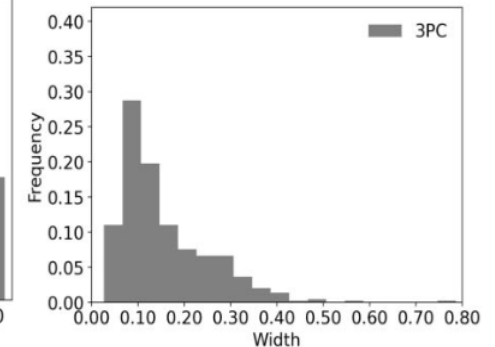
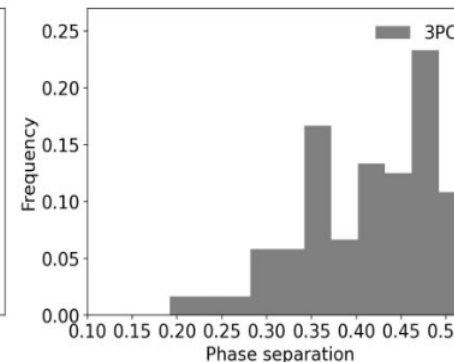
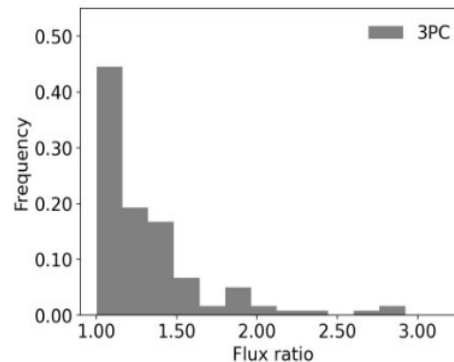
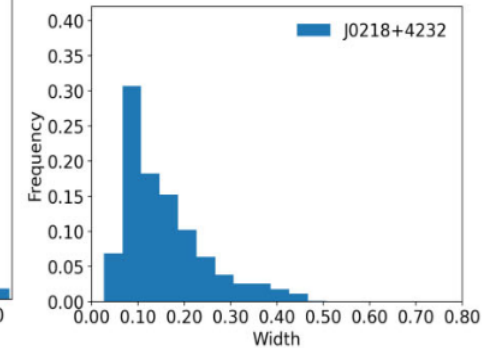
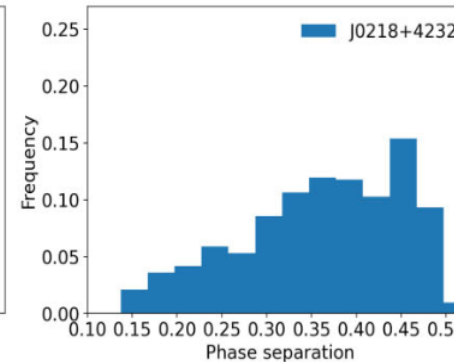
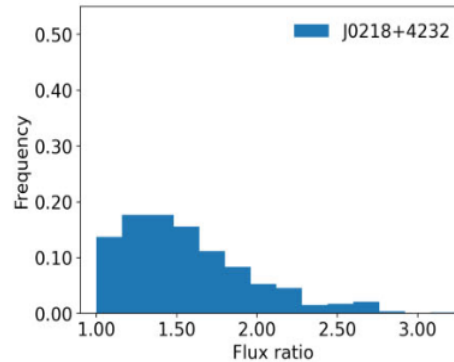
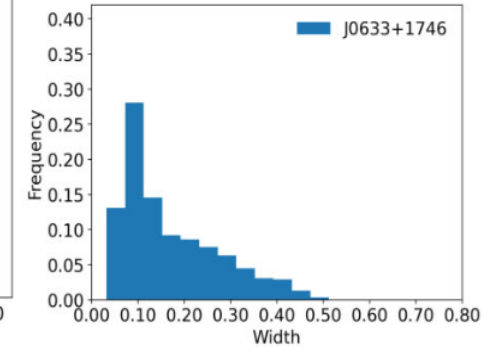
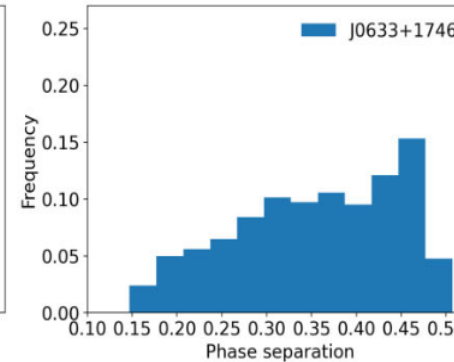
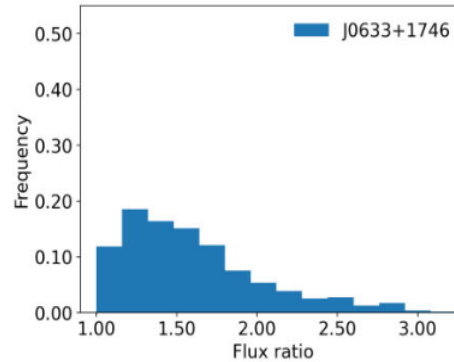
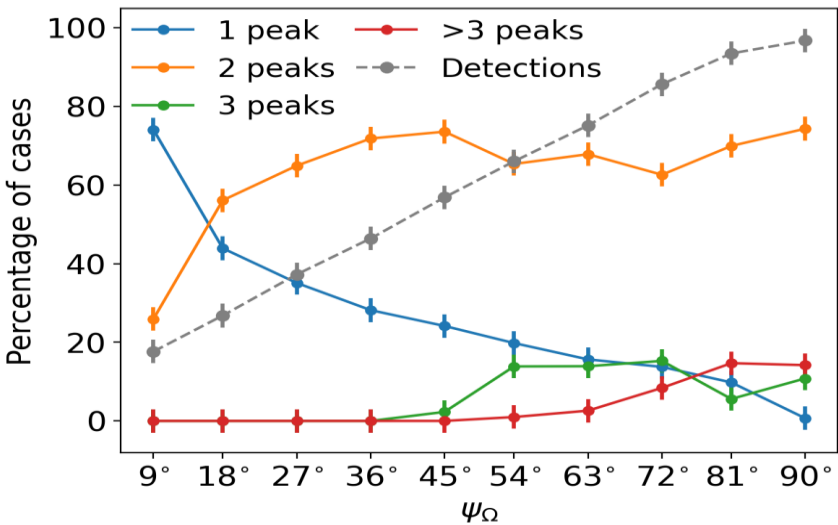
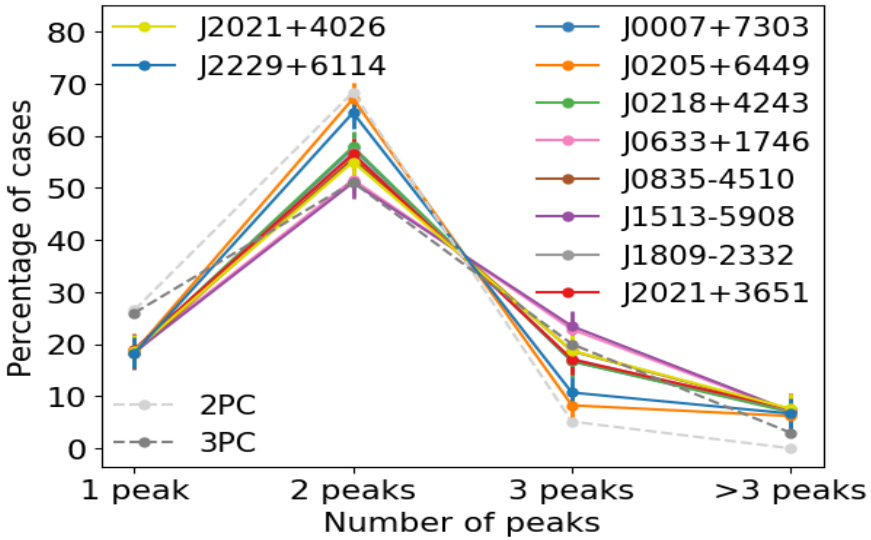
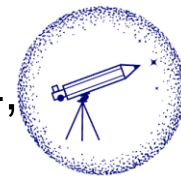
- We build emission maps, from which light curves are obtained



[Íñiguez-Pascual,
Torres & Viganò
2024, MNRAS,
530, 1550]

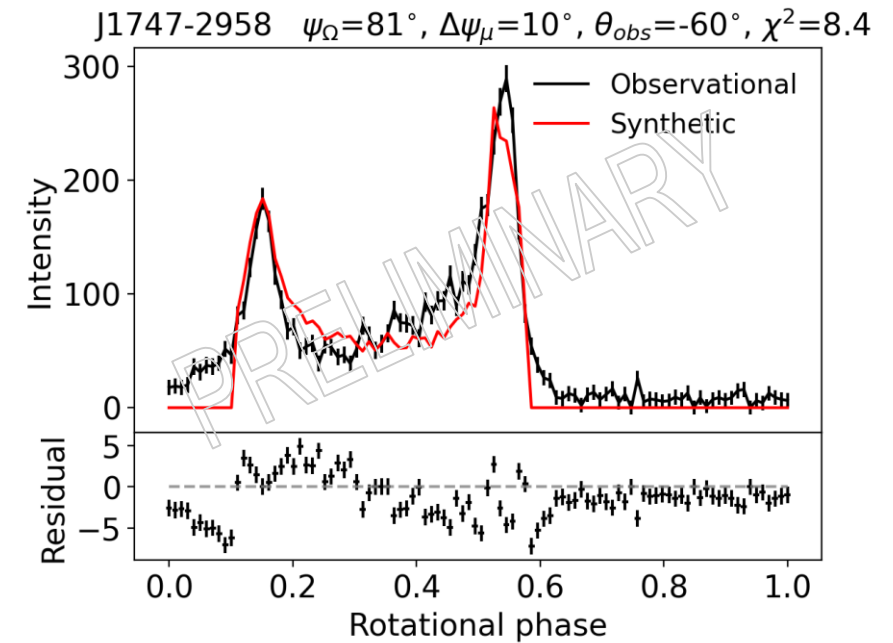
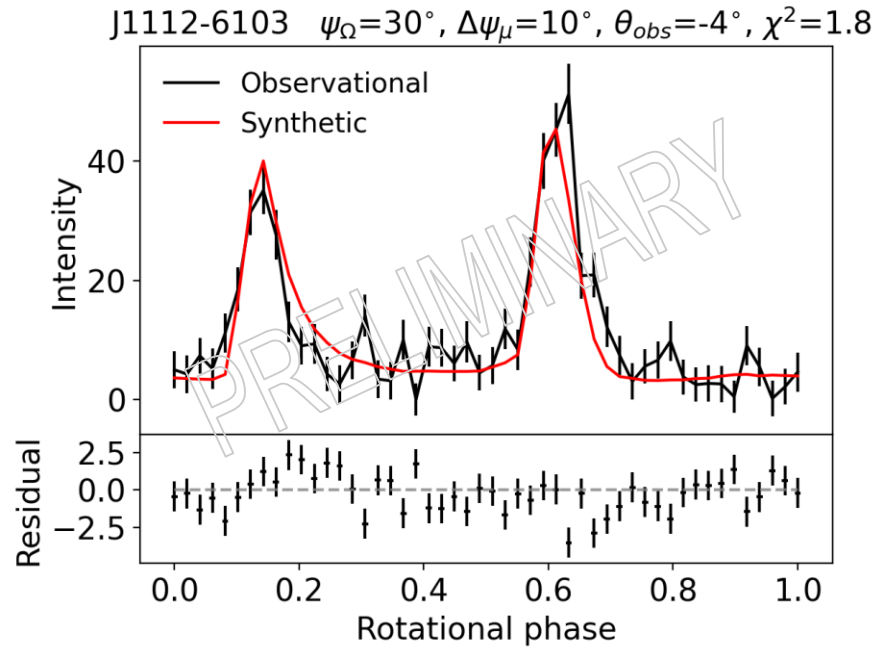
Light curves statistics

[Íñiguez-Pascual, Torres & Viganò 2024, MNRAS, 530, 1550]



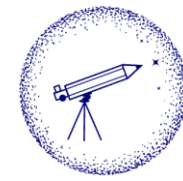
Light curves fitting

- Fitting synthetic light curves to observational ones, concurrently to the spectral fitting

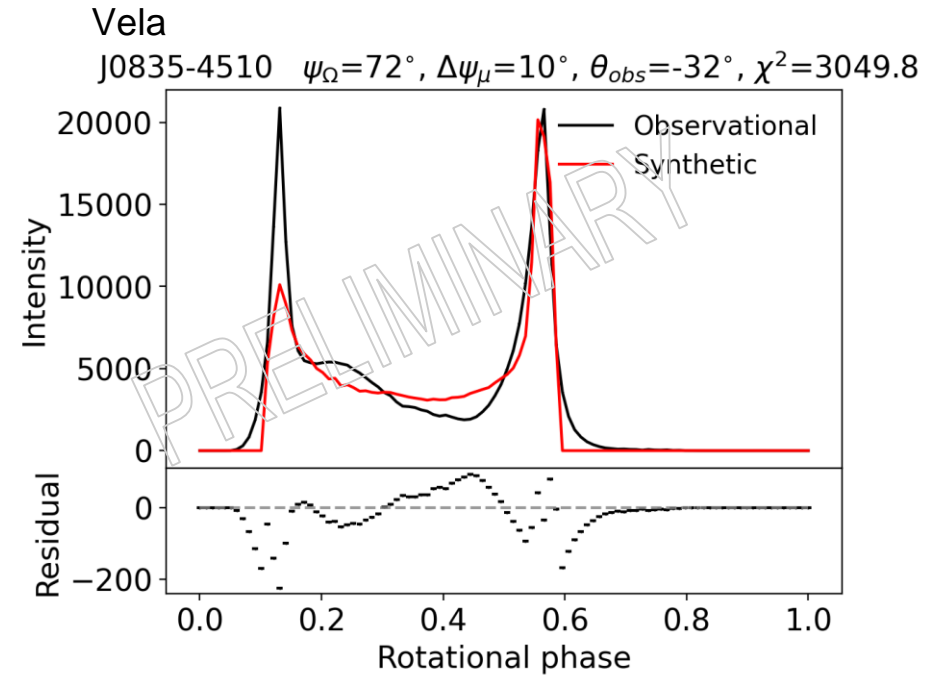
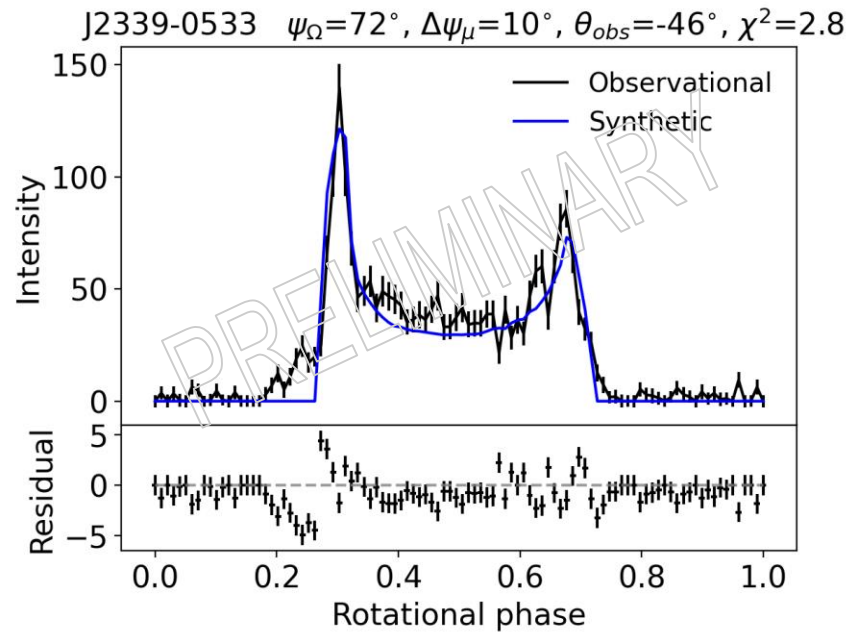


[Íñiguez-Pascual
et al. (in prep.)]

Light curves fitting

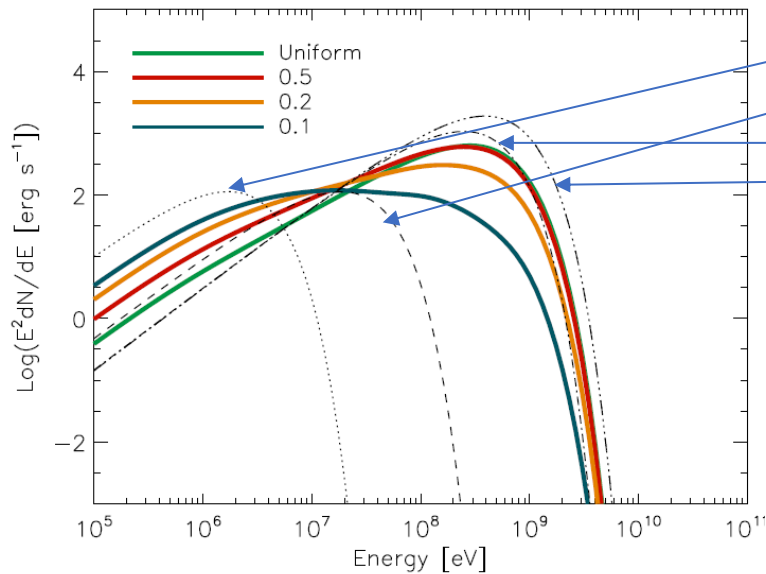
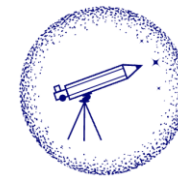


- Fitting synthetic light curves to observational ones, concurrently to the spectral fitting



[Íñiguez-Pascual
et al. (in prep.)]

Caveats



contributions at:
 $r/R_{lc} = 0.01$
 $r/R_{lc} = 0.1$
 $r/R_{lc} = 0.5$
 $r/R_{lc} = 1$

$$\frac{dN}{dx} = N_0 \frac{e^{-(x-x_{min})/x_0}}{x_0(1 - e^{-(x_{max}-x_{min})/x_0})}$$

Without this "weight", the spectra would be basically the single-particle spectra (uniform), which have a fixed slope and cannot fit most of the spectra. The weight could be physically:

- A geometrical effect -> Discarded, we still need it when we compute the emission map
- The result of neglecting complications: variation of electric field along a line and across the accelerating region, presence of backward particles, better prescription for the curvature radius...

A result of this is that the spectra and gamma-ray light curves are substantially decoupled: the gamma-ray light curves are determined by the geometry, with little dependence on the spectra.

Conclusions

- Our spectral model is able to properly fit the population of high-energy emitting pulsars, showing that synchro-curvature radiation is an appropriate mechanism to explain the emission from these objects

[Íñiguez-Pascual D., Viganò D., Torres D. F., 2022, MNRAS, 516, 2475 (2208.05549)]

- The geometrical model reproduces the variety of observational gamma-ray light curves (though, by definition, it cannot capture their small scale features)

[Íñiguez-Pascual D., Torres D. F., Viganò D., 2024, MNRAS, 530, 1550 (2404.01926)]

Future prospects

- Soon to be submitted: full-sample concurrent spectral and light curves fitting
- Include more realistic physics to improve the major caveats (arbitrary weights in the spectral model, simple recipes for electric field and curvature radius), while keeping our effective approach

Thank you



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